

Spin and Orbital Rotation of Electrons and Photons via Spin-Orbit Interaction

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INTRODUCTION

- Photons and electrons have 4 degrees of freedom
- In cylindrical symmetry, these may be taken as:
 - Spectral (frequency): ω
 - Polarization (spin, or SAM): σ
 - Orbital angular momentum (OAM): m_ℓ (z component)
 - “Radial” quantum, number: n

INTRODUCTION

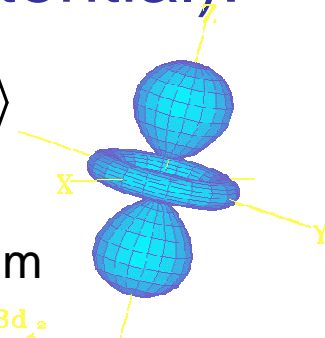
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Hydrogen atom
(Spherical potential):

Cylindrical potential: $\Psi \propto e^{i(\beta z - \omega t)}$

$\ell \rightarrow \omega, \beta$

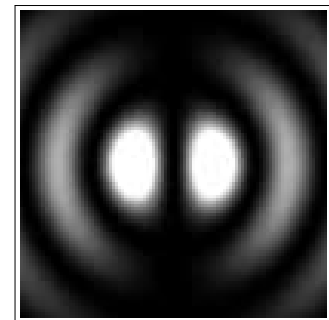
$$|n \ell m_\ell \sigma\rangle$$



$\hat{\mathbf{L}} \rightarrow \ell$ (Good quantum number)
3d_{z²}

$$|\omega m_\ell \sigma n\rangle$$

\updownarrow
 β



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ω

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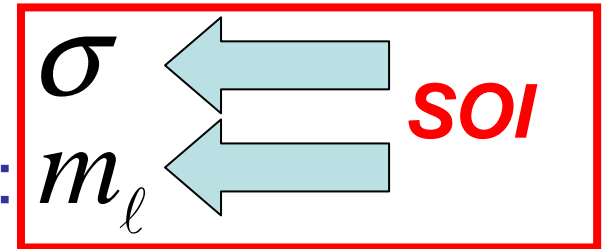
σ

- Orbital angular momentum (OAM):

m_ℓ

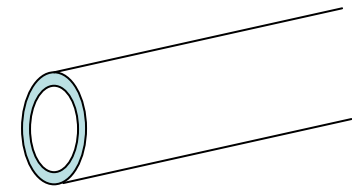
- “Radial” quantum, number:

n



- All previously studied (most recently photon OAM)
- **This talk concerns:**
 - The interaction between the SAM and OAM degrees of freedom of electrons and photons in cylindrical symmetry
 - The transfer of entanglement between the SAM and OAM degrees of freedom for two particle states

OUTLINE (elevator pitch)



1: Electron in an inhomogeneous potential– effective \mathbf{B} field induces spin-orbit interaction

2: A photon in an inhomogeneous medium experiences an analogous spin-orbit interaction, even though it has no magnetic moment

3: In a cylindrical waveguide, the spin-orbit dynamics are described by a single expression applying to both electrons and photons

4: This leads to the prediction of several spin (SAM) and orbital (OAM) rotational effects:

- a) spatial or time evolution of either particle's spin/polarization vector is controlled by the sign of its OAM
- b) conversely, its spatial wavefunction evolution is controlled by its SAM

5: Two-particle states: Reversible transfer of entanglement between SAM and OAM

Question: Why such a close analogy?

quant-ph > arXiv:0905.3778

To what extent is there a photon-electron analogy? An Example:

1. Spin Hall Effect for Electrons: opposite spin accumulation on opposing lateral surfaces of a current-carrying sample. Its origin is spin-orbit interaction.

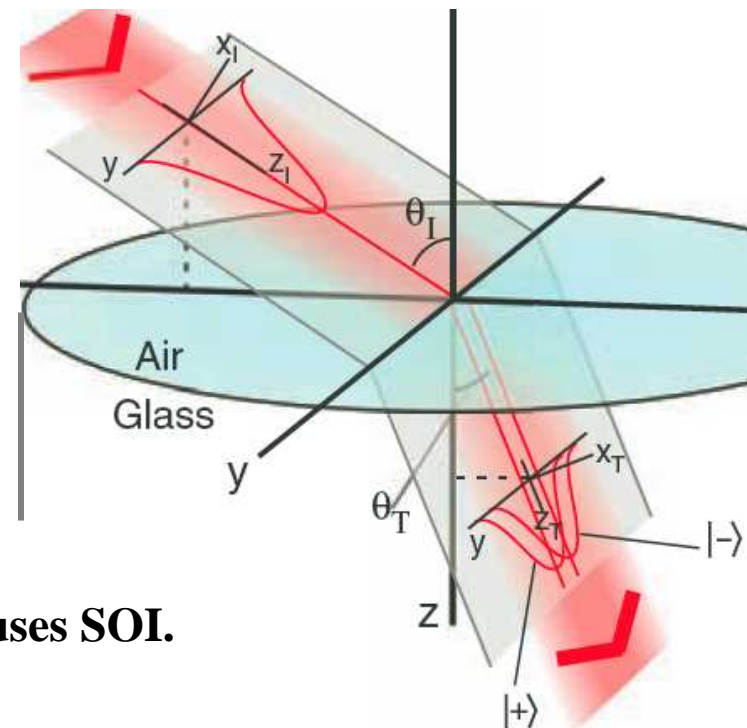
Dyakonov and Perel (1971) Sov. Phys. JETP Lett. 13, 467

Hirsch (1999) PRL 83, 1834

2. Spin Hall Effect for Light: spin-dependent displacement perpendicular to the refractive index gradient for photons passing through an air-glass interface.

M. Onoda, S. Murakami, N. Nagaosa, PRL 93, 083901 (2004)

Observed: Hosten, Kwiat Science 319 (2008)



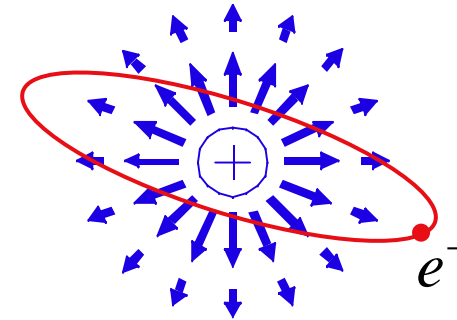
Inhomogeneity in refractive index causes SOI.

Spin-Orbit Interaction (SOI) in Spherical Potentials:

- **ELECTRON** IN AN INHOMOGENOUS SPHERICAL ELECTRIC POTENTIAL (ATOM)

$$H' = -\frac{e^2}{2m^2c^2} \frac{1}{r} \frac{\partial V}{\partial r} \mathbf{S} \cdot \mathbf{L}$$

$$\mathbf{S} = SAM \quad \mathbf{r} \times \mathbf{p} = \mathbf{L} = OAM$$



(atomic fine structure)

Spin-Orbit Interaction (SOI) in Cylindrical Potentials?

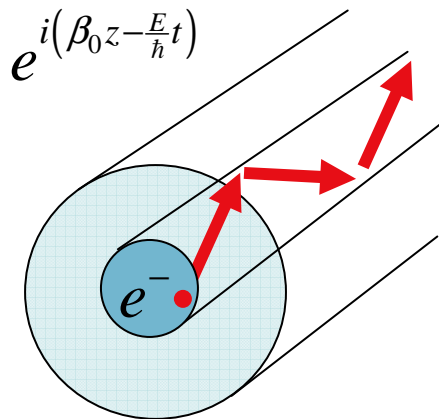
C Leary, D Reeb, M Raymer, NJP, 10, 103022 (2008).

$$H' = -\frac{e^2}{2m^2c^2} \frac{1}{\rho} \frac{\partial V}{\partial \rho} S_z L_z$$

$$S_z = \mathbf{S} \cdot \hat{\mathbf{z}}$$

$$L_z = \mathbf{L} \cdot \hat{\mathbf{z}}$$

$$\Psi \propto e^{i(\beta_0 z - \frac{E}{\hbar} t)}$$



Traveling wave

Schrodinger equations: $\hat{H}\Psi = E\Psi$
Energy shift

$\hat{H}\Psi = \beta^2\Psi$
Propagation constant shift
(phase velocity)

Perhaps the cylindrical electron case is not so surprising...

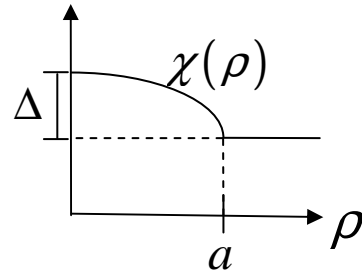
$$\hat{H}\Psi = \beta^2\Psi$$

$$\Psi_0 \propto e^{i(\beta_0 z - \frac{E}{\hbar}t)}$$

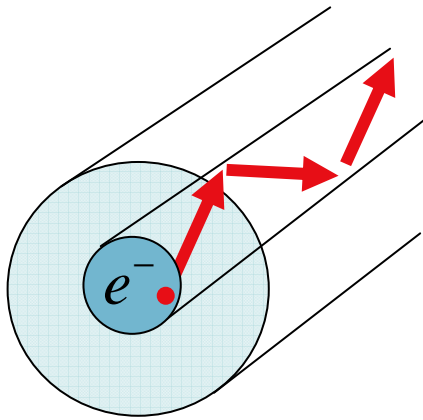
$$\hat{H} = \hat{H}_0 + \hat{H}'$$

Solve Schrodinger
(Dirac) equation:

$$V(\rho) \propto -U(\rho)$$



$$U(\rho) = (U(0) - \Delta\chi(\rho))$$



$$\hat{H}'_{e^-} = -\frac{\Delta}{4\beta_0} \frac{1}{\rho} \frac{\partial \chi}{\partial \rho} S_z L_z$$

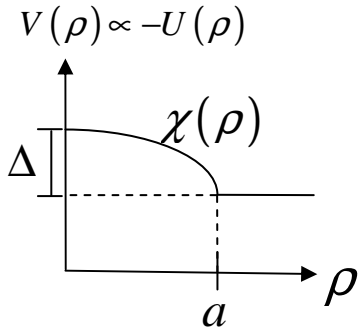
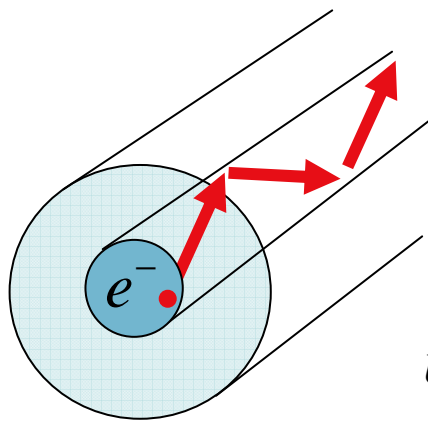
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$$\hat{H}\Psi = \beta^2\Psi$$

$$\Psi_0 \propto e^{i(\beta_0 z - \frac{E}{\hbar}t)}$$

$$\hat{H} = \hat{H}_0 + \hat{H}'$$

Solve Schrodinger
(Dirac) equation:



$$U(\rho) = (U(0) - \Delta\chi(\rho))$$

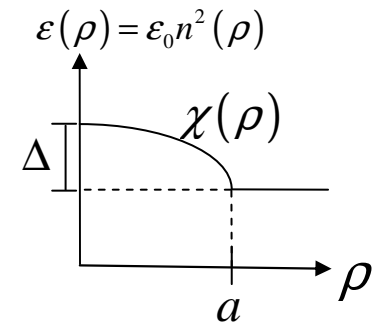
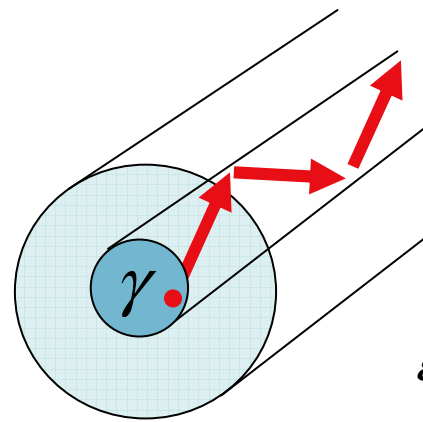
$$\hat{H}'_{e^-} = -\frac{\Delta}{4\beta_0} \frac{1}{\rho} \frac{\partial \chi}{\partial \rho} S_z L_z$$

$$\hat{H}E = \beta^2 E$$

$$\Psi_0 \propto e^{i(\beta_0 z - \omega t)}$$

$$\hat{H} = \hat{H}_0 + \hat{H}'$$

Solve Helmholtz
(Maxwell) equations:



$$\varepsilon(\rho) = \varepsilon(0)(1 - \Delta\chi(\rho))$$

$$\hat{H}'_{\gamma} = -\frac{\Delta}{4\beta_0} \frac{1}{\rho} \frac{\partial \chi}{\partial \rho} S_z L_z$$

but it is remarkable that the photon SOI is completely analogous!

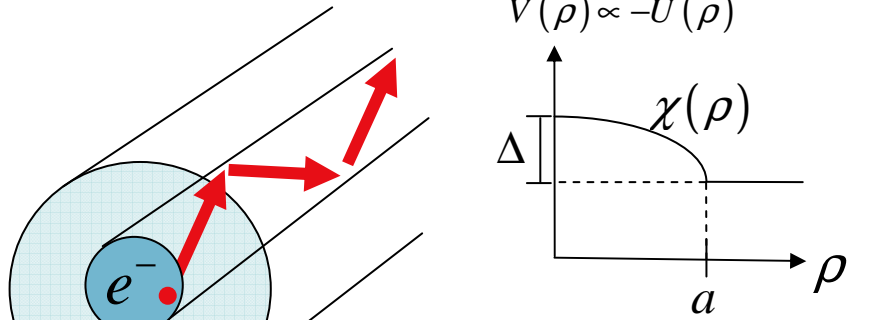
Perhaps the cylindrical electron case is not so surprising...

$$\hat{H}\Psi = \beta^2\Psi$$

$$\Psi_0 \propto e^{i(\beta_0 z - \frac{E}{\hbar}t)}$$

Solve Schrodinger (Dirac) equation:

$V(\rho) \propto -U(\rho)$



$U(\rho) = (U(0) - \Delta\chi(\rho))$

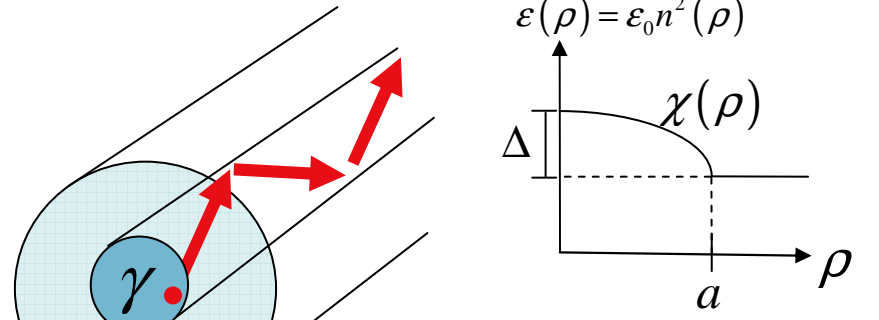
$$\hat{H}'_{e^-} = -\frac{\Delta}{4\beta_0} \frac{1}{\rho} \frac{\partial \chi}{\partial \rho} S_z L_z$$

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Solve Helmholtz (Maxwell) equations:

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$\varepsilon(\rho) = \varepsilon(0)(1 - \Delta\chi(\rho))$

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but it is remarkable that the photon SOI is completely analogous!

Main Point: the $S_z L_z$ part of the Hamiltonian means that β will undergo a small positive/negative shift depending upon whether S_z and L_z are oriented parallel/antiparallel to each other

Wavefunction: $\Psi_{\sigma m_\ell} = \psi(\rho) e^{-im_\ell \phi} e^{i(\beta_0 z - \omega t)} \hat{\mathbf{e}}_\sigma$

\uparrow OAM \uparrow SAM

This slide common to both particle types

$m_\ell = 0, \pm 1, \pm 2, \dots =$ OAM quantum number $\sigma = \pm 1 =$ SAM quantum number

Hamiltonian: $H'_{SOI} = -\frac{\Delta}{4\beta_0} \frac{1}{\rho} \frac{\partial \chi}{\partial \rho} S_z L_z \propto S_z L_z$

Perturbation theory: $\delta\beta = \langle \Psi_{\sigma m_\ell} | H'_{SOI} | \Psi_{\sigma m_\ell} \rangle \propto \sigma m_\ell \Rightarrow \Psi_{\sigma m_\ell} \rightarrow \Psi_{\sigma m_\ell} e^{i\delta\beta z}$

“Notation”: $\Psi_{+1,+2} = \psi(\rho) e^{+2i\phi} e^{i(\beta_0 z - \omega t)} \hat{\mathbf{e}}_{+1} \Leftrightarrow \text{⊕} \text{ etc. (no more math!)}$

rotation effect: $\text{⊕} + \text{⊖} = \text{⊕↔}$

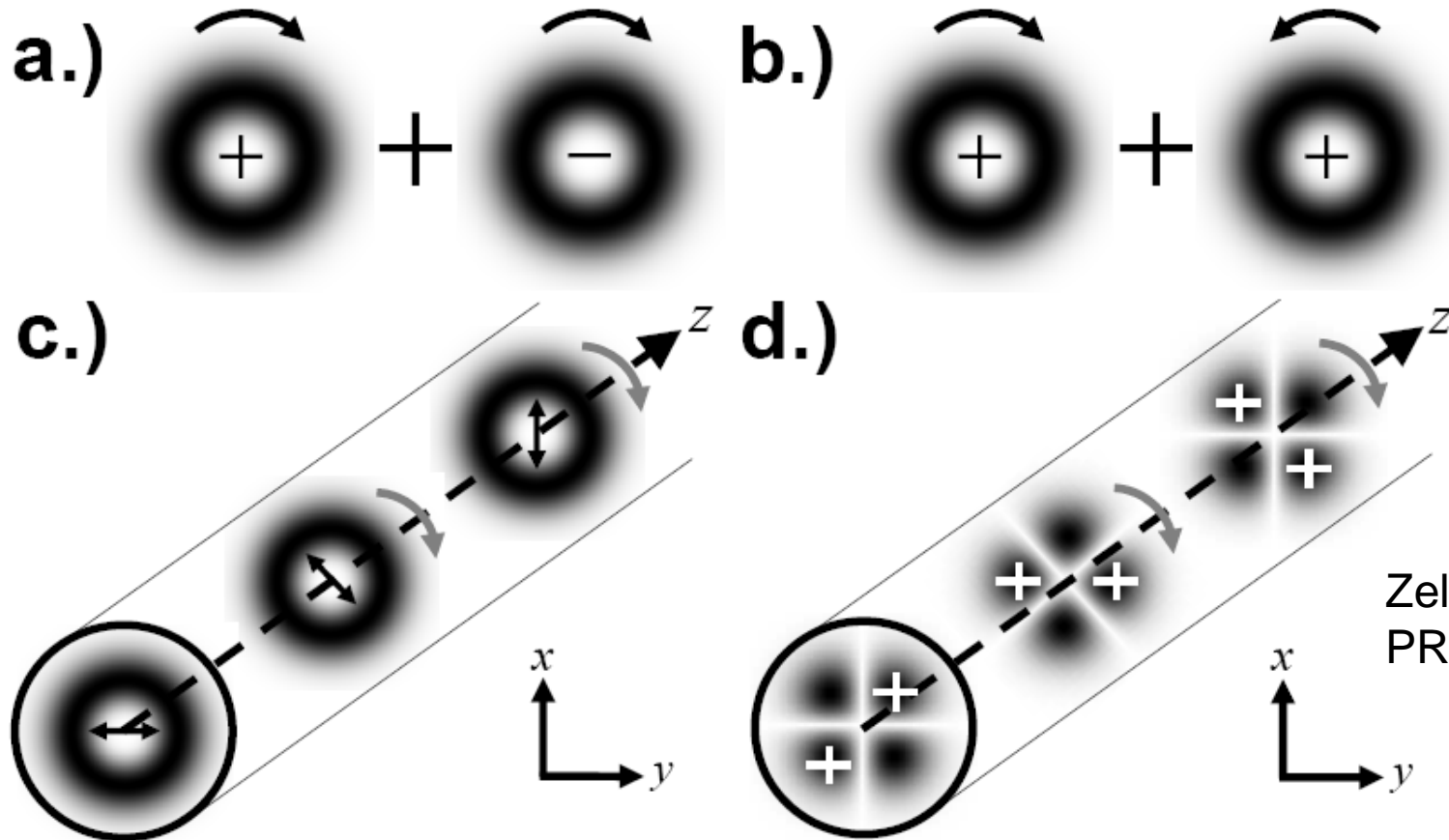
$\rightarrow \text{⊕} e^{+i|\delta\beta|z} + \text{⊖} e^{-i|\delta\beta|z}$

$\propto \cos(|\delta\beta|z) \hat{\mathbf{x}} + \mu \sin(|\delta\beta|z) \hat{\mathbf{y}}$

Spin/polarization rotation direction controlled by sign of $\mu = \pm 1 \equiv \frac{m_\ell}{|m_\ell|}$

“orbit”-controlled “spin” rotation:

Complementary spin-orbit effects



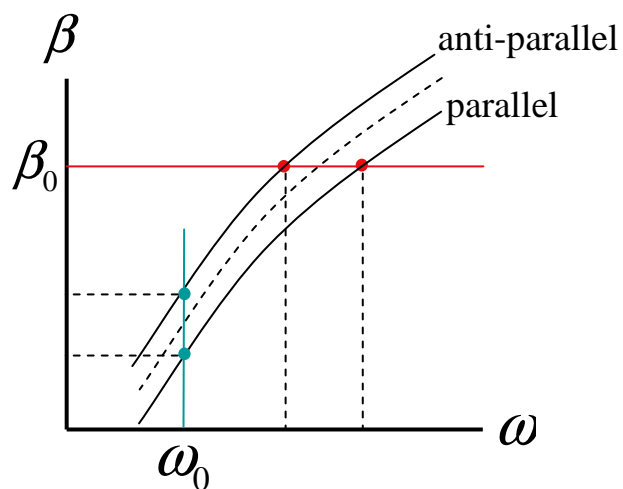
Zeldovich,
PRA '91

“orbit”-controlled “spin”
rotation (OAM eigenstate)

“spin”-controlled “orbit”
rotation (SAM eigenstate)

Both of these effects may occur either in space or time

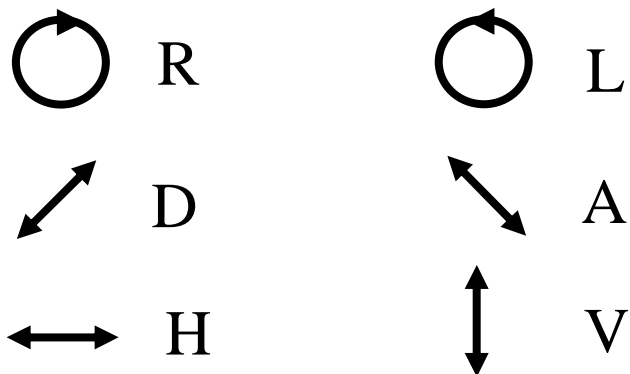
Spatial vs. Temporal Rotation



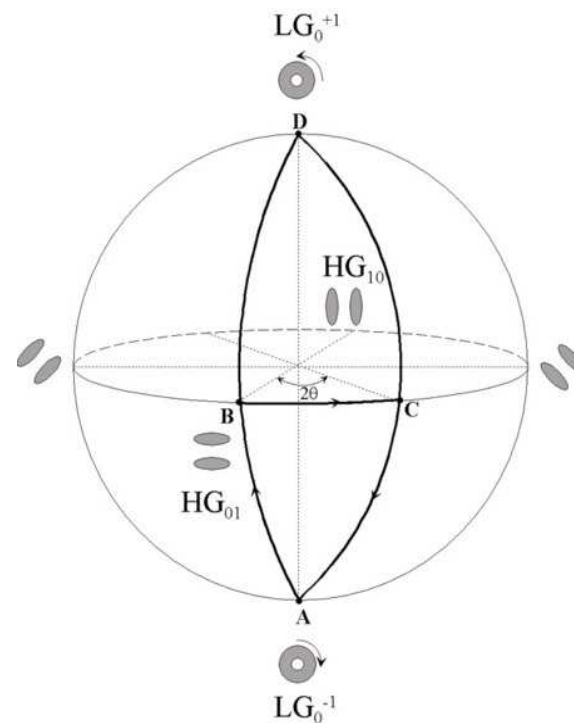
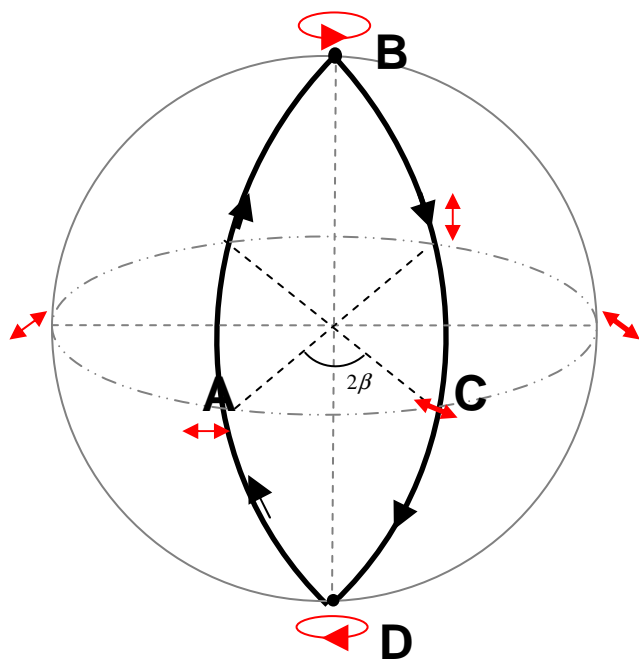
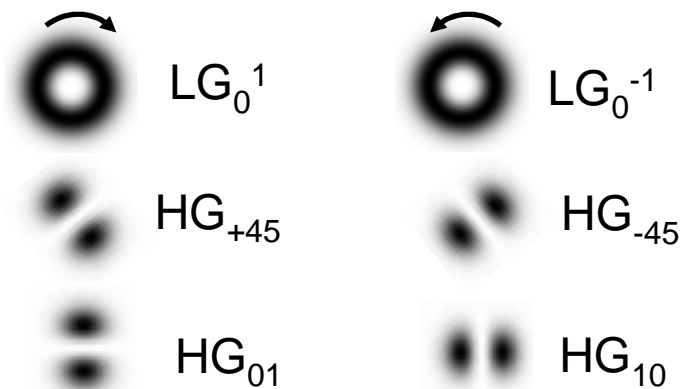
- The SOI may be thought of as a splitting of the dispersion curve for “parallel” vs. “anti-parallel” states.
- For a given frequency, there are two β values, so that the SOI splits the propagation constant (spatial rotation)
- If both the frequencies and propagation constants of the parallel and anti-parallel states are originally slightly different, then the SOI acts to *restore* the propagation constants to degeneracy (temporal rotation)

SAM to OAM entanglement transfer

Polarization (SAM)

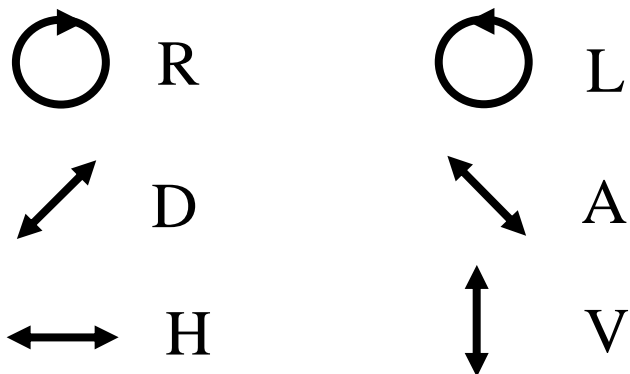


Transverse Beam Shape (OAM)

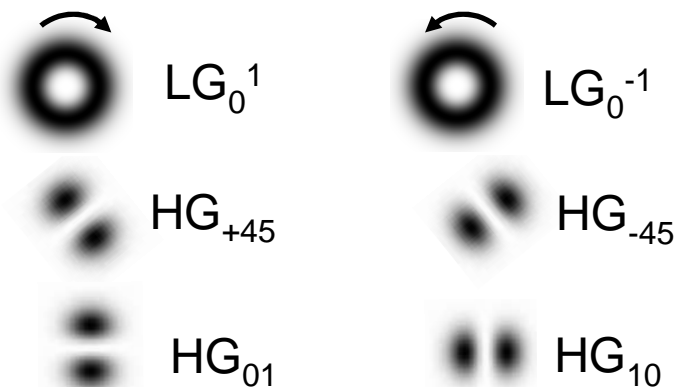


SAM to OAM entanglement transfer

Polarization (SAM)



Transverse Beam Shape (OAM)



Begin in a purely polarization Bell state:

$$|H \text{ " " } \rangle |V \text{ " " } \rangle + |V \text{ " " } \rangle |H \text{ " " } \rangle$$

Quarter wave plates:

$$|R \text{ " " } \rangle |L \text{ " " } \rangle + |L \text{ " " } \rangle |R \text{ " " } \rangle$$

“Orbital” rotation SOI gate:

$$|R \text{ " " } \rangle |L \text{ " " } \rangle + |L \text{ " " } \rangle |R \text{ " " } \rangle$$

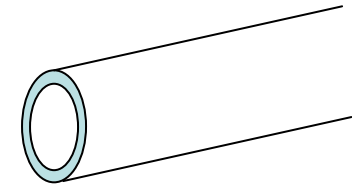
Quarter wave plates
and mode converters:

$$|D \text{ " " } \rangle |A \text{ " " } \rangle + |A \text{ " " } \rangle |D \text{ " " } \rangle$$

“Spin” rotation SOI gate:
(end in purely OAM-entangled state)

$$|H \text{ " " } \rangle |H \text{ " " } \rangle + |H \text{ " " } \rangle |H \text{ " " } \rangle$$

OUTLINE REVISITED

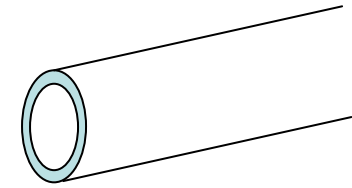


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- 3: In a cylindrical waveguide, the spin-orbit dynamics are described by a single expression applying to both electrons and photons
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Question: Why such a close analogy?

quant-ph > arXiv:0905.3778

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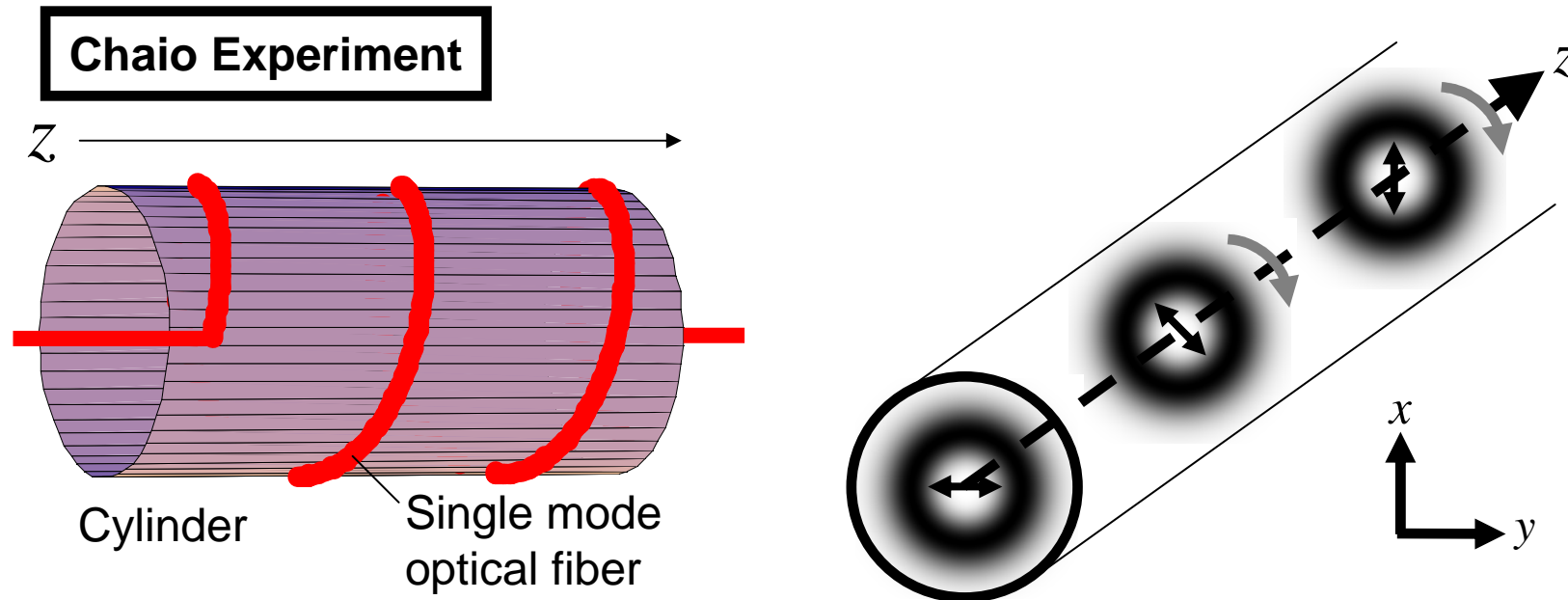
Question: Why such a close analogy?

Answer: The spin-orbit interaction is a purely geometric effect (Berry phase)

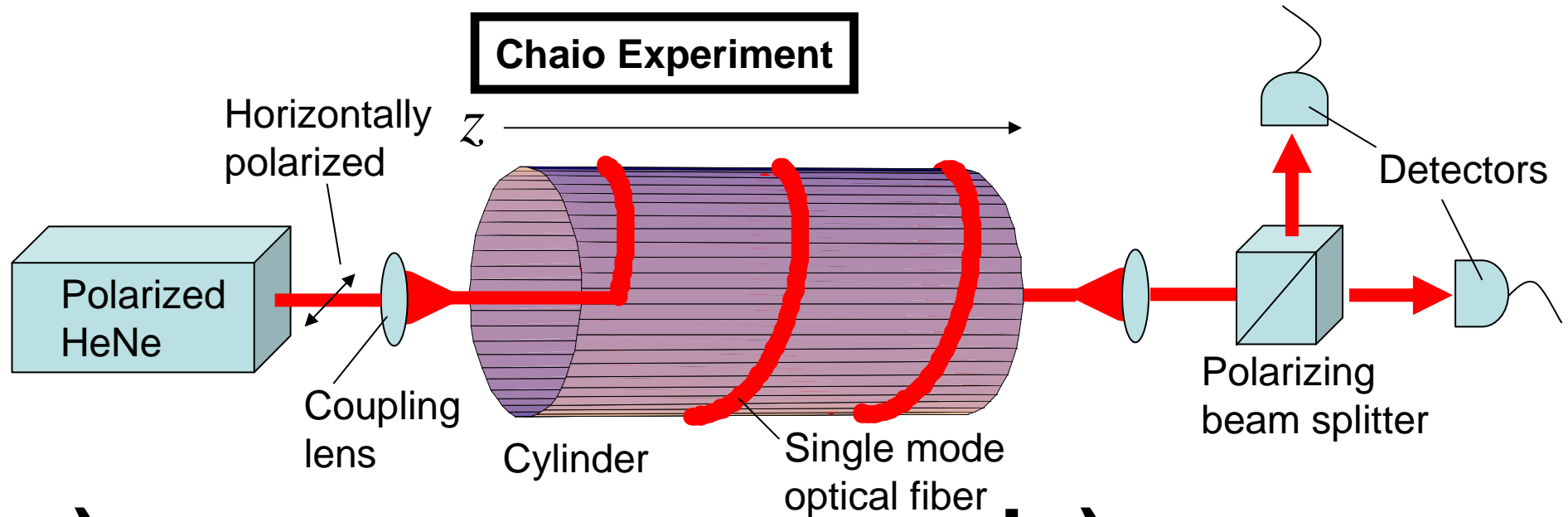
quant-ph > arXiv:0905.3778

Seem familiar?

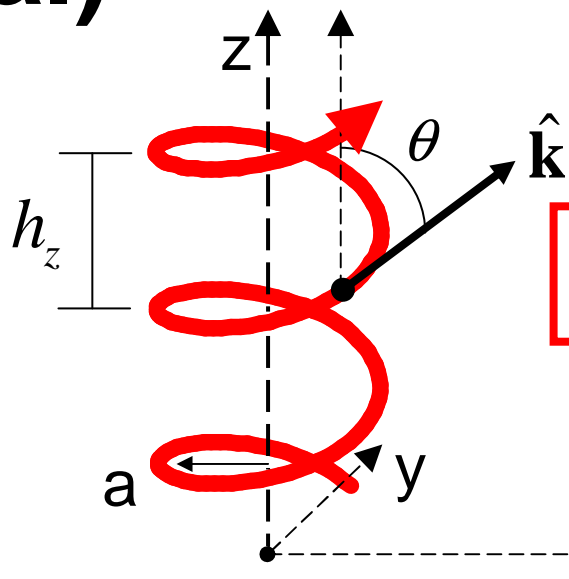
The polarization rotation effect is reminiscent of an experiment of Chaio, in which the geometric (Berry) phase was first observed as polarization rotation in a coiled fiber.



However, the Chaio effect involves only the fundamental Gaussian mode, which has zero OAM, and predicts an accumulated geometric phase of zero for a straight fiber! Nevertheless, considering this effect provides a way in which to understand the straight fiber rotation phenomena for modes with OAM.



a.)

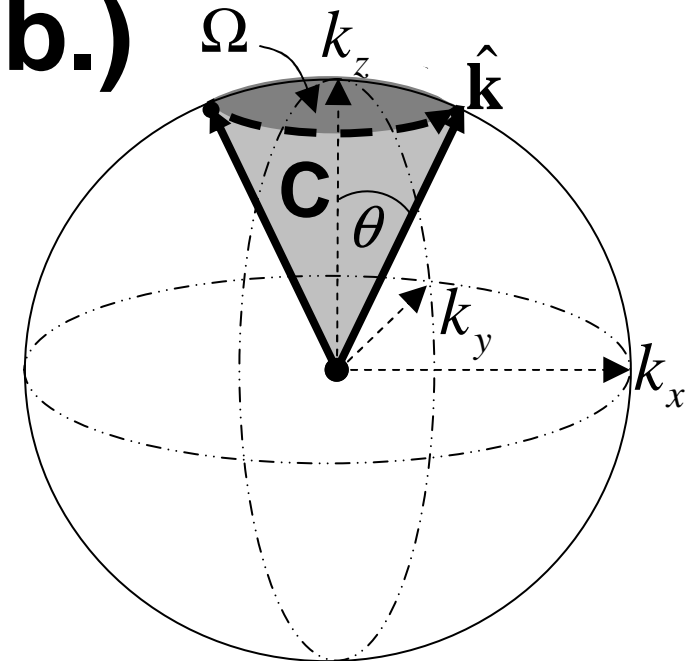


$$\Omega = 2\pi N (1 - \cos \theta)$$

Where $N = \#$ of coils

$$E \rightarrow E e^{-i\sigma\mu_{cl}\Omega}$$

b.)



The geometric phase equals the surface area enclosed by the curve C , which is traced out by the momentum vector \mathbf{k} . This phase causes a “beating” effect: polarization rotation

Geometric phase approach

- “Classical” geometric phase:
 (“Classical” particle with spin; Bialynicki-Birula)

$$E \rightarrow E e^{-i\sigma\mu_{cl}\Omega(\mathbf{k})} \equiv E e^{-i\gamma}$$

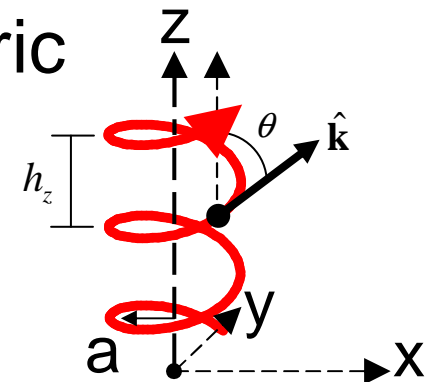
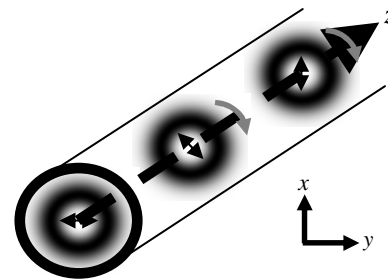
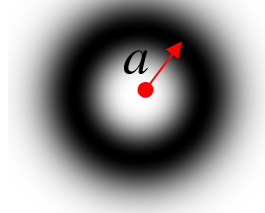
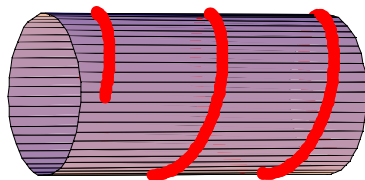
- For quantum states, we quantize:

$$\gamma \rightarrow \hat{\gamma} \equiv -\sigma\mu \frac{\Delta}{2\beta_0 a^2} \nabla_T$$

- The geometric phase $\gamma = \delta\beta z$ is then given by the expectation of Ω in the unperturbed states:

$$\gamma = \delta\beta z = \langle \Psi_{\sigma m_\ell} | \hat{\gamma} | \Psi_{\sigma m_\ell} \rangle \propto \sigma m_\ell \Rightarrow \Psi_{\sigma m_\ell} \rightarrow \Psi_{\sigma m_\ell} e^{i\gamma} = \Psi_{\sigma m_\ell} e^{i\delta\beta z}$$

- If a is set equal to the radius of the wavefunction “peak”, then the perturbative and geometric approaches agree well.

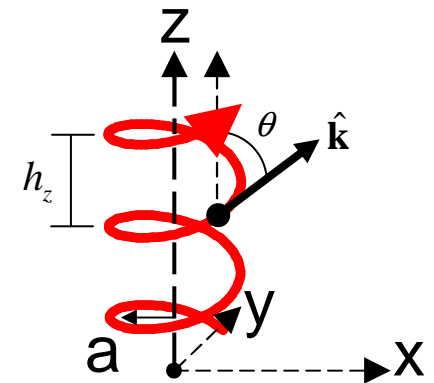
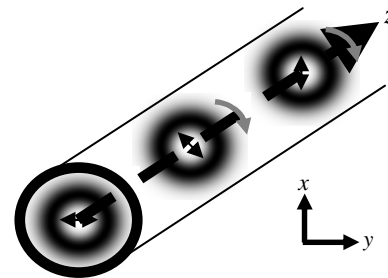
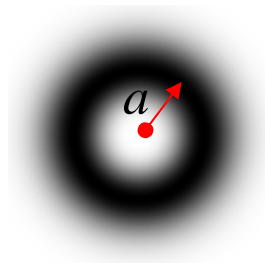
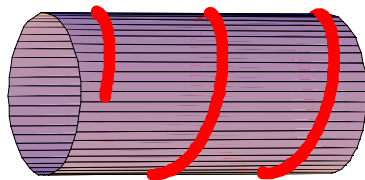


Geometric phase approach

- Specifically, one gets a particularly simple form for the accumulated phase $e^{i\delta\beta z}$ in the special case of a step-index/potential of large radius a , for OAM states which are “near cutoff” – that is, their wavefunction “peak” is near the fiber radius:

$$\gamma = \delta\beta z = \sigma m_\ell \frac{\Delta}{2\beta_0 a^2} z$$

- In this case the two approaches agree exactly.



Conclusions

- The spin-orbit interaction dynamics of the electron and photon are identical to first order in perturbation theory
- They have a common geometric origin
- The role of the electron's potential energy is played by the permittivity in the photon case
- Any particle with spin (i.e. neutrons) will also undergo SOI
- The SOI allows for the construction of both spin and orbital “gates”, which may enable reversible transfer of entanglement between SAM and OAM
- Experiments are underway to observe these effects